

# Désordre et phénomènes critiques : les modèles d'accrochage

Giambattista Giacomin

Université Paris Diderot and Laboratoire Probabilités et Modèles Aléatoires (LPMA)

30.08.2016

# Motivations

## General issue

- Much work to establish existence of phase transitions for "statistical mechanics systems" and to understand the behavior of these systems approaching the phase transition (critical phenomena: role of exactly solvable models)

# Motivations

## General issue

- Much work to establish existence of phase transitions for "statistical mechanics systems" and to understand the behavior of these systems approaching the phase transition (critical phenomena: role of exactly solvable models)
- But what is the effect of disorder on critical phenomena? Lots of physical predictions. . .

# Motivations

## General issue

- Much work to establish existence of phase transitions for "statistical mechanics systems" and to understand the behavior of these systems approaching the phase transition (critical phenomena: role of exactly solvable models)
- But what is the effect of disorder on critical phenomena? Lots of physical predictions. . .

Polymer pinning is a class of statistical mechanics models

- with an exactly solvable character (without disorder)

# Motivations

## General issue

- Much work to establish existence of phase transitions for "statistical mechanics systems" and to understand the behavior of these systems approaching the phase transition (critical phenomena: role of exactly solvable models)
- But what is the effect of disorder on critical phenomena? Lots of physical predictions. . .

Polymer pinning is a class of statistical mechanics models

- with an exactly solvable character (without disorder)
- for which much has been done when disorder is introduced

# Motivations

## General issue

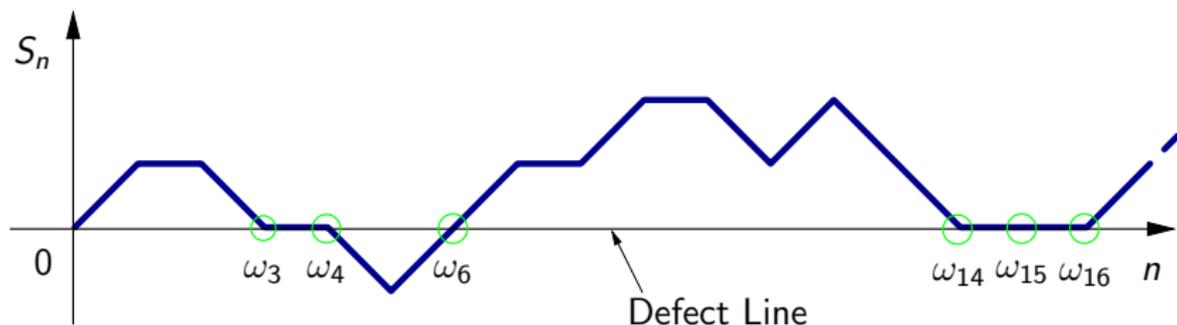
- Much work to establish existence of phase transitions for "statistical mechanics systems" and to understand the behavior of these systems approaching the phase transition (critical phenomena: role of exactly solvable models)
- But what is the effect of disorder on critical phenomena? Lots of physical predictions. . .

Polymer pinning is a class of statistical mechanics models

- with an exactly solvable character (without disorder)
- for which much has been done when disorder is introduced
- interesting from a modeling viewpoint

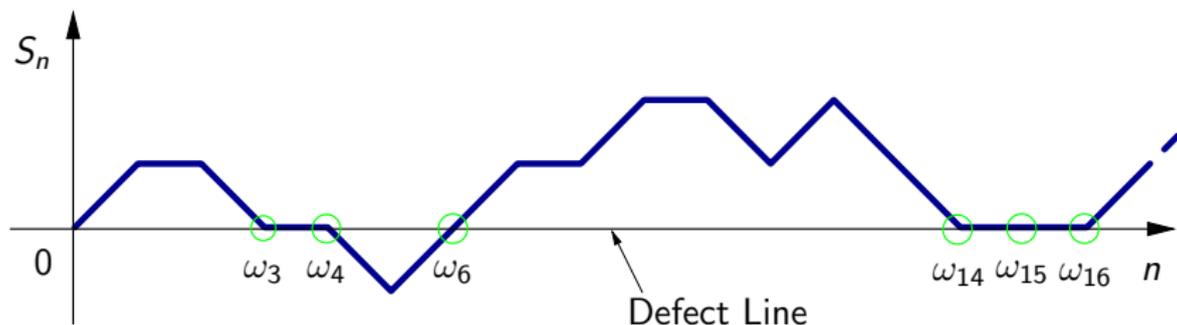
# “The Example”: $(1 + d)$ -directed walk models

Symmetric Random Walk  $\{S_n\}_n$  with increments in  $\{-1, 0, +1\}$  ( $d = 1$ )



# “The Example”: $(1 + d)$ -directed walk models

Symmetric Random Walk  $\{S_n\}_n$  with increments in  $\{-1, 0, +1\}$  ( $d = 1$ )



Model ( $\beta \geq 0$ ,  $h \in \mathbb{R}$ ):

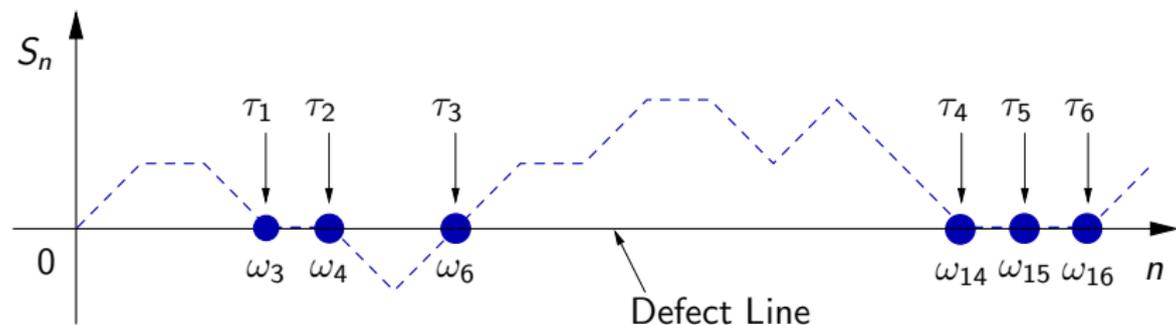
$$\frac{d\mathbf{P}_{N,\omega}}{d\mathbf{P}}(S) = \frac{1}{Z_{N,\omega}} \exp\left(\sum_{n=1}^N (\beta\omega_n + h) \delta_n\right)$$

with  $\delta_n = \mathbf{1}_{S_n=0}$ . The disorder  $\omega$  is a IID sequence  $\mathcal{N}(0, 1)$  of law  $\mathbb{P}$ .



# 'Rethinking 'The Example''

Symmetric Random Walk  $\{S_n\}_n$  with increments in  $\{-1, 0, +1\}$



Model ( $\beta \geq 0, h \in \mathbb{R}$ ):

$$\frac{d\mathbf{P}_{N,\omega}}{d\mathbf{P}}(\tau) = \frac{1}{Z_{N,\omega}} \exp\left(\sum_{n=1}^N (\beta\omega_n + h) \delta_n\right)$$

with  $\delta_n = \mathbf{1}_{n \in \tau}$ .

# In the end: renewal based model

Basic building block: a Discrete Renewal Process

# In the end: renewal based model

Basic building block: a Discrete Renewal Process

- 1  $\tau = \{\tau_0, \tau_1, \tau_2, \dots\}$  discrete renewal sequence (that is,  $\tau_0 = 0$  and  $\{\tau_j - \tau_{j-1}\}_{j \in \mathbb{N}}$  is IID), of law  $\mathbf{P}$ , s. t.

$$K(n) = \mathbf{P}(\tau_1 = n) \sim C_K/n^{1+\alpha}, \quad (C_K > 0),$$

and

$$\sum_{n \in \mathbb{N}} K(n) \leq 1.$$

# In the end: renewal based model

Basic building block: a Discrete Renewal Process

- ①  $\tau = \{\tau_0, \tau_1, \tau_2, \dots\}$  discrete renewal sequence (that is,  $\tau_0 = 0$  and  $\{\tau_j - \tau_{j-1}\}_{j \in \mathbb{N}}$  is IID), of law  $\mathbf{P}$ , s. t.

$$K(n) = \mathbf{P}(\tau_1 = n) \sim C_K/n^{1+\alpha}, \quad (C_K > 0),$$

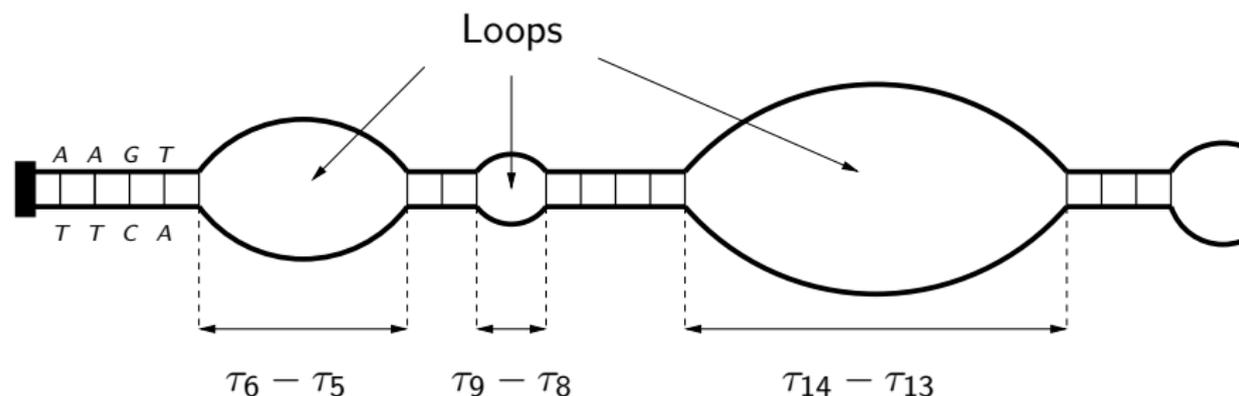
and

$$\sum_{n \in \mathbb{N}} K(n) \leq 1.$$

- ② If  $\sum_n K(n) < 1 \implies$  renewal on  $\mathbb{N} \cup \{\infty\}$ , with  $K(\infty) = 1 - \sum_n K(n)$  (*terminating renewal*), otherwise the renewal is *persistent*.

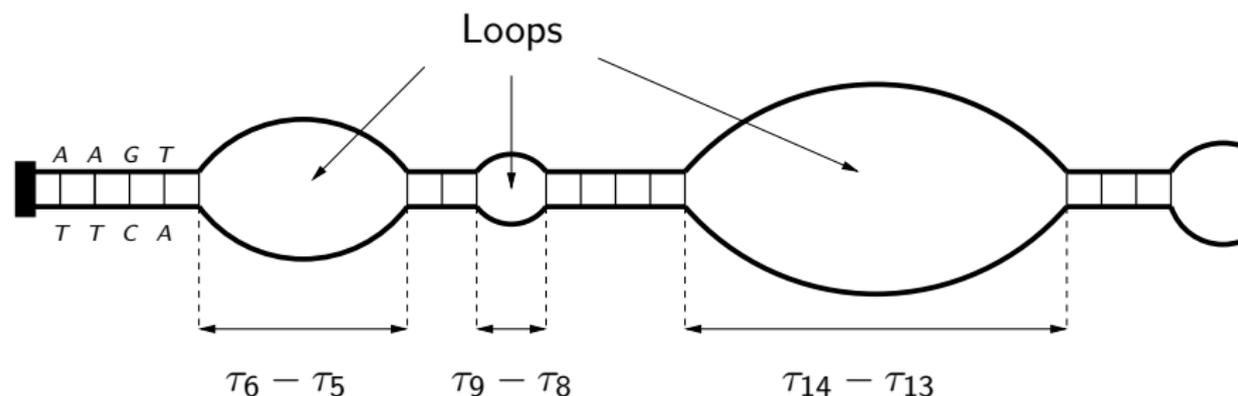
Obs. :  $\alpha = 1/2$  for both  $d = 1$  and  $3$ , but  $\sum_n K(n) < 1$  if  $d = 3$

# The Poland-Scheraga model (DNA denaturation)



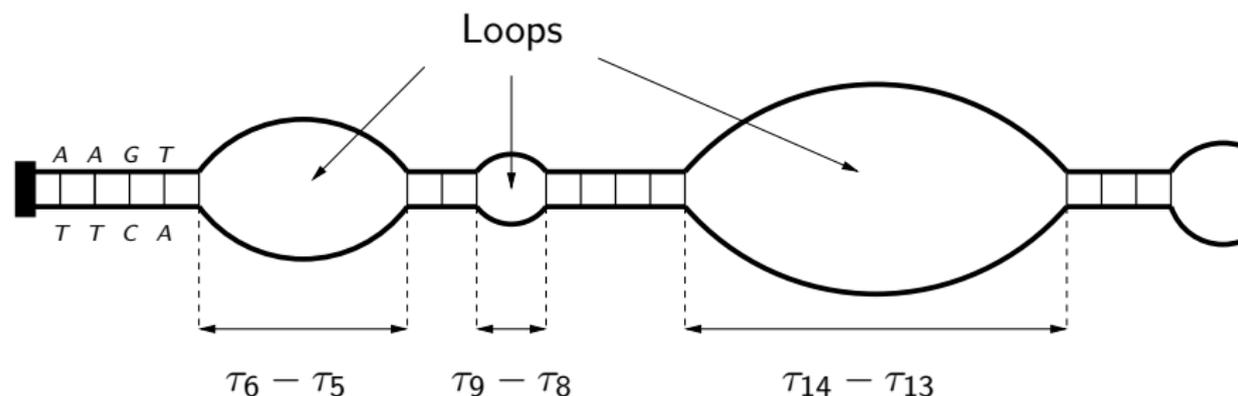
- The two thick lines are the DNA strands. They may be paired, gaining thus energetic contributions that depend on whether the base pair is A-T or G-C.
- There are then sections of unpaired bases (the *loops*) to which an entropy is associated: loops correspond to inter-arrival of length  $n \geq 2$ .

# The Poland-Scheraga model (DNA denaturation)



- The two thick lines are the DNA strands. They may be paired, gaining thus energetic contributions that depend on whether the base pair is A-T or G-C.
- There are then sections of unpaired bases (the *loops*) to which an entropy is associated: loops correspond to inter-arrival of length  $n \geq 2$ .
- Entropy:  $C\mu^n/n^{1+\alpha}$ ,  $\alpha \approx 1.15$  and  $\mu$  drops out when  $N \rightarrow \infty$ .

# The Poland-Scheraga model (DNA denaturation)



- The two thick lines are the DNA strands. They may be paired, gaining thus energetic contributions that depend on whether the base pair is A-T or G-C.
- There are then sections of unpaired bases (the *loops*) to which an entropy is associated: loops correspond to inter-arrival of length  $n \geq 2$ .
- Entropy:  $C\mu^n/n^{1+\alpha}$ ,  $\alpha \approx 1.15$  and  $\mu$  drops out when  $N \rightarrow \infty$ .

## Annealed, pure and homogeneous models

$$Z_{N,\omega} = \mathbf{E} \left[ \exp \left( \sum_{n=1}^N (\beta\omega_n + h) \delta_n \right) \right]$$

so

$$\mathbb{E} Z_{N,\omega,\beta,h} = \mathbf{E} \left[ \exp \left( \sum_{n=1}^N ((\beta^2/2) + h) \delta_n \right) \right]$$

which is the partition function of a homogeneous model.

## Annealed, pure and homogeneous models

$$Z_{N,\omega} = \mathbf{E} \left[ \exp \left( \sum_{n=1}^N (\beta\omega_n + h) \delta_n \right) \right]$$

so

$$\mathbb{E} Z_{N,\omega,\beta,h} = \mathbf{E} \left[ \exp \left( \sum_{n=1}^N ((\beta^2/2) + h) \delta_n \right) \right]$$

which is the partition function of a homogeneous model.

So:

- The annealed (or *pure*) model is just a homogeneous model with pinning potential  $h + \beta^2/2$ ;

## Annealed, pure and homogeneous models

$$Z_{N,\omega} = \mathbf{E} \left[ \exp \left( \sum_{n=1}^N (\beta\omega_n + h) \delta_n \right) \right]$$

so

$$\mathbb{E} Z_{N,\omega,\beta,h} = \mathbf{E} \left[ \exp \left( \sum_{n=1}^N ((\beta^2/2) + h) \delta_n \right) \right]$$

which is the partition function of a homogeneous model.

So:

- The annealed (or *pure*) model is just a homogeneous model with pinning potential  $h + \beta^2/2$ ;
- Homogeneous pinning models are *exactly solvable*

# The (quenched) free energy (density)

Fundamental expression:

$$\lim_{N \rightarrow \infty} \frac{1}{N} \mathbb{E} \log Z_{N,\omega} =: F(\beta, h)$$

# The (quenched) free energy (density)

Fundamental expression:

$$\lim_{N \rightarrow \infty} \frac{1}{N} \mathbb{E} \log Z_{N,\omega} =: F(\beta, h)$$

Path properties (more natural questions for probabilists?) are subordinated to understanding the free energy behavior.

# The (quenched) free energy (density)

Fundamental expression:

$$\lim_{N \rightarrow \infty} \frac{1}{N} \mathbb{E} \log Z_{N,\omega} =: F(\beta, h)$$

Path properties (more natural questions for probabilists?) are subordinated to understanding the free energy behavior.

Important facts

- $F(\cdot, \cdot)$  is convex,  $F(\beta, \cdot)$  and  $F(\cdot, h)$  are non-decreasing
- Jensen's inequality yields

$$\mathbb{E} \log Z_{N,\omega} \leq \log \mathbb{E} Z_{N,\omega}$$

# The (quenched) free energy (density)

Fundamental expression:

$$\lim_{N \rightarrow \infty} \frac{1}{N} \mathbb{E} \log Z_{N,\omega} =: F(\beta, h)$$

Path properties (more natural questions for probabilists?) are subordinated to understanding the free energy behavior.

Important facts

- $F(\cdot, \cdot)$  is convex,  $F(\beta, \cdot)$  and  $F(\cdot, h)$  are non-decreasing
- Jensen's inequality yields

$$\mathbb{E} \log Z_{N,\omega} \leq \log \mathbb{E} Z_{N,\omega}$$

so

$$F(0, h) \leq F(\beta, h) \leq F(0, h + \beta^2/2)$$

# The (quenched) free energy (density)

Fundamental expression:

$$\lim_{N \rightarrow \infty} \frac{1}{N} \mathbb{E} \log Z_{N,\omega} =: F(\beta, h)$$

Path properties (more natural questions for probabilists?) are subordinated to understanding the free energy behavior.

Important facts

- $F(\cdot, \cdot)$  is convex,  $F(\beta, \cdot)$  and  $F(\cdot, h)$  are non-decreasing
- Jensen's inequality yields

$$\mathbb{E} \log Z_{N,\omega} \leq \log \mathbb{E} Z_{N,\omega}$$

so

$$F(0, h) \leq F(\beta, h) \leq F(0, h + \beta^2/2)$$

- $F(\beta, h) \geq 0$ , so either  $= 0$  or  $> 0$

# Localization and Delocalization

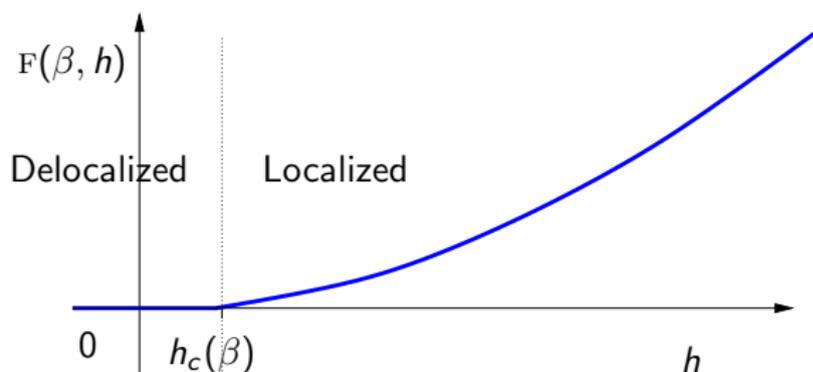
Proof of  $F(\beta, h) \geq 0$ :

$$\begin{aligned} F(\beta, h) &= \lim_N \frac{1}{N} \mathbb{E} \log \mathbf{E} \left[ \exp \left( \sum_{n=1}^N (\beta \omega_n + h) \delta_n \right) \right] \\ &\geq \lim_N \inf \frac{1}{N} \mathbb{E} \log \mathbf{E} \left[ \exp \left( \sum_{n=1}^N (\beta \omega_n + h) \delta_n \right) ; \tau_1 > N \right] \\ &= \lim_N \frac{1}{N} \log \mathbf{P} (\tau_1 > N) = 0. \quad \square \end{aligned}$$

# Localization and Delocalization

Proof of  $F(\beta, h) \geq 0$ :

$$\begin{aligned} F(\beta, h) &= \lim_N \frac{1}{N} \mathbb{E} \log \mathbf{E} \left[ \exp \left( \sum_{n=1}^N (\beta \omega_n + h) \delta_n \right) \right] \\ &\geq \lim_N \inf \frac{1}{N} \mathbb{E} \log \mathbf{E} \left[ \exp \left( \sum_{n=1}^N (\beta \omega_n + h) \delta_n \right) ; \tau_1 > N \right] \\ &= \lim_N \frac{1}{N} \log \mathbf{P} (\tau_1 > N) = 0. \quad \square \end{aligned}$$



# Main questions

## Plenty of questions, but above all:

- Can one compute or estimate  $h_c(\beta)$ ?
- Critical behavior?  $F(\beta, h) \underset{h \searrow h_c(\beta)}{\sim} \text{const.}(h - h_c(\beta))^{\nu_q}$
- [Of course also: path properties]

# Main questions

## Plenty of questions, but above all:

- Can one compute or estimate  $h_c(\beta)$ ?
- Critical behavior?  $F(\beta, h) \underset{h \searrow h_c(\beta)}{\sim} \text{const.} (h - h_c(\beta))^{\nu_q}$
- [Of course also: path properties]

The pure (homogeneous, annealed) model is solvable

# Main questions

## Plenty of questions, but above all:

- Can one compute or estimate  $h_c(\beta)$ ?
- Critical behavior?  $F(\beta, h) \stackrel{h \searrow h_c(\beta)}{\sim} \text{const.}(h - h_c(\beta))^{\nu_q}$
- [Of course also: path properties]

The pure (homogeneous, annealed) model is solvable

Semi-explicit formula for  $F(0, h)$  from which

- $h_c(0) = -\log \sum_n K(n) (\geq 0)$



$$F(0, h) \stackrel{h \searrow h_c(0)}{\sim} \text{const.}(h - h_c(0))^{\nu_a}, \quad \nu_a = \max(1/\alpha, 1)$$

...M. Fisher '84. But: Erdos, Pollard, Feller, Garsia, Lamperti... (40's...)

# General principles to deal with disorder(?)

## Recall the main questions:

- Compute or estimate  $h_c(\beta)$
- Critical behavior?  $F(\beta, h) \underset{h \searrow h_c(\beta)}{\sim} \text{const.}(h - h_c(\beta))^{\nu_q}$

# General principles to deal with disorder(?)

## Recall the main questions:

- Compute or estimate  $h_c(\beta)$
- Critical behavior?  $F(\beta, h) \underset{h \searrow h_c(\beta)}{\sim} \text{const.}(h - h_c(\beta))^{\nu_q}$

## Harris Criterion (A. B. Harris 1974)

Knowing the critical behavior of the pure system one can decide whether (at small disorder) the critical behavior of pure and disordered systems coincide (the disorder is irrelevant) or differ (the disorder is relevant).

# General principles to deal with disorder(?)

## Recall the main questions:

- Compute or estimate  $h_c(\beta)$
- Critical behavior?  $F(\beta, h) \stackrel{h \searrow h_c(\beta)}{\sim} \text{const.}(h - h_c(\beta))^{\nu_q}$

## Harris Criterion (A. B. Harris 1974)

Knowing the critical behavior of the pure system one can decide whether (at small disorder) the critical behavior of pure and disordered systems coincide (the disorder is irrelevant) or differ (the disorder is relevant).

HC for pinning models [Forgacs et al. (1986), Derrida et al. (1992)]:

- $h_c(\beta) = h_c^a(\beta)$  and  $\nu_q = \nu_a$  for  $\beta \leq \beta_0$  if  $\alpha < 1/2$
- $h_c(\beta) > h_c^a(\beta)$  and (**probably**)  $\nu_q \neq \nu_a$  for  $\beta > 0$  and  $\alpha > 1/2$ .
- marginal case (controversial issue for physicists)

# Harris criterion for pinning models: rigorous results

- Full understanding of "non-controversial" physical predictions (with sharper estimates), both in the irrelevant and relevant disorder regime

# Harris criterion for pinning models: rigorous results

- Full understanding of "non-controversial" physical predictions (with sharper estimates), both in the irrelevant and relevant disorder regime
- Solution of controversial marginal case (precise estimates on  $h_c(\beta)$ )

# Harris criterion for pinning models: rigorous results

- Full understanding of "non-controversial" physical predictions (with sharper estimates), both in the irrelevant and relevant disorder regime
- Solution of controversial marginal case (precise estimates on  $h_c(\beta)$ )
- Proof of  $\nu_q > \nu_a$  for  $\alpha > 1/2$ , but (**main open problem**)  $\nu_q$  remains unknown

# Harris criterion for pinning models: rigorous results

- Full understanding of "non-controversial" physical predictions (with sharper estimates), both in the irrelevant and relevant disorder regime
- Solution of controversial marginal case (precise estimates on  $h_c(\beta)$ )
- Proof of  $\nu_q > \nu_a$  for  $\alpha > 1/2$ , but (**main open problem**)  $\nu_q$  remains unknown

## Key specificity for pinning model

Results have gone so far because one of the two phases (the delocalized one) is trivial from the free energy viewpoint

# Harris criterion for pinning models: rigorous results

- Full understanding of "non-controversial" physical predictions (with sharper estimates), both in the irrelevant and relevant disorder regime
- Solution of controversial marginal case (precise estimates on  $h_c(\beta)$ )
- Proof of  $\nu_q > \nu_a$  for  $\alpha > 1/2$ , but (**main open problem**)  $\nu_q$  remains unknown

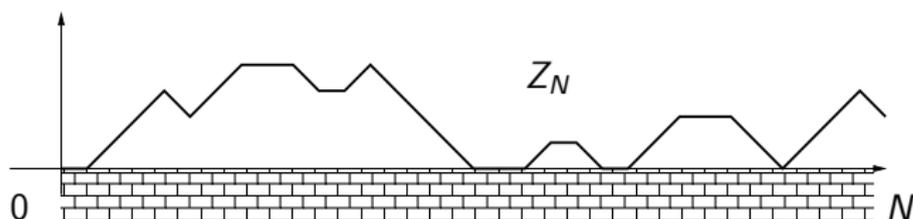
## Key specificity for pinning model

Results have gone so far because one of the two phases (the delocalized one) is trivial from the free energy viewpoint

[Alexander, Berger, Biskup, Bolthausen, Caravenna, Derrida, G., den Hollander, Lacoïn, Opoku, Pétrélis, Poisat, Sohier, Sun, Toninelli, Torri, Zygouras: 2004-ongoing]

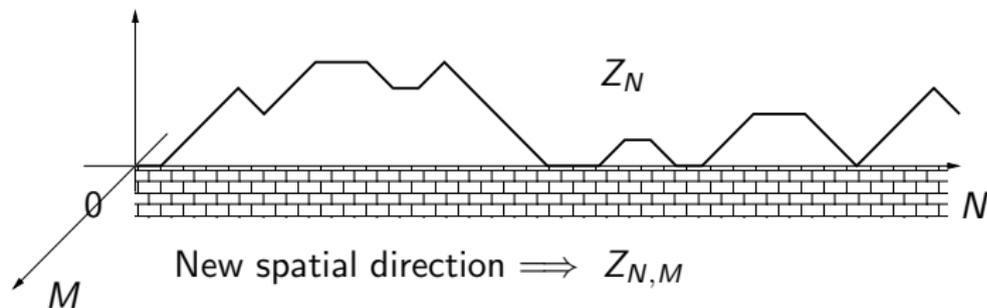
## Recent developments: higher dimensional

- Pinning of  $d + 1$  dimensional interfaces (free field, SOS type models) [Coquille, Milos], [G., Lacoïn]



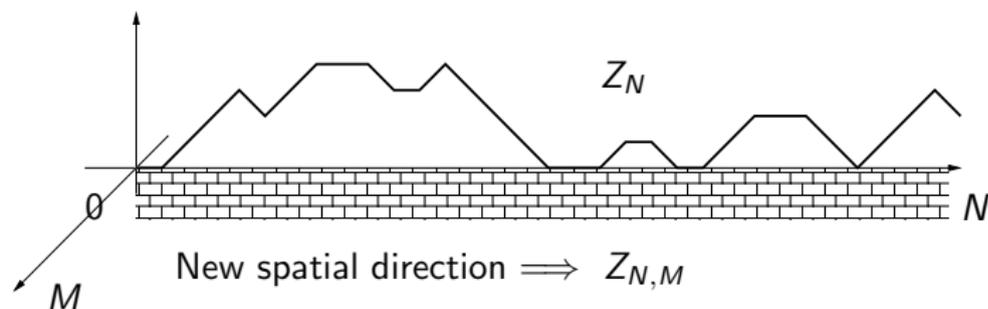
# Recent developments: higher dimensional

- Pinning of  $d + 1$  dimensional interfaces (free field, SOS type models) [Coquille, Milos], [G., Lacoïn]



# Recent developments: higher dimensional

- Pinning of  $d + 1$  dimensional interfaces (free field, SOS type models) [Coquille, Milos], [G., Lacoïn]



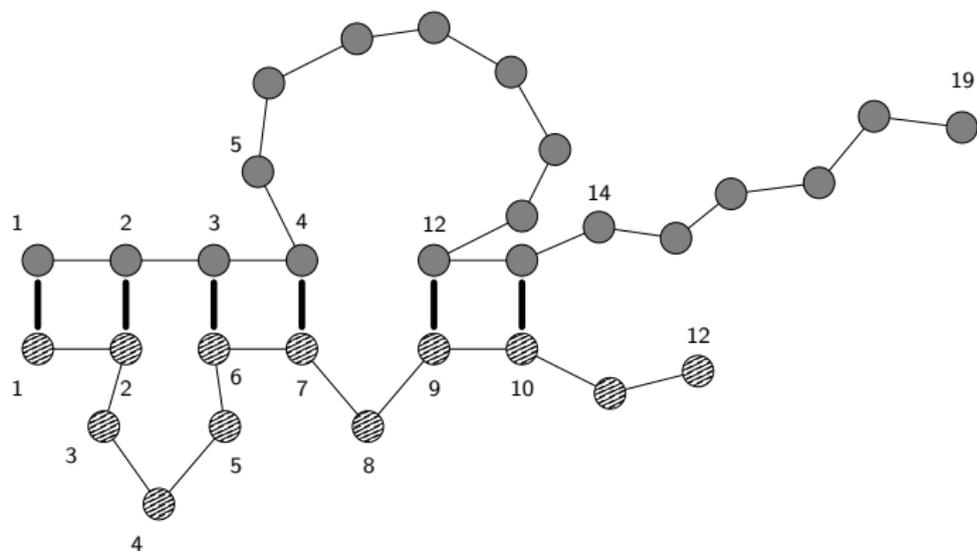
- Pinning of multi-dimensional renewal: (multi  $\rightarrow$  two)

$$\frac{d\mathbf{P}_{N,M,\omega}}{d\mathbf{P}}(\tau) = \frac{1}{Z_{N,M,\omega}} \exp \left( \sum_{n=1}^N \sum_{m=1}^M (\beta\omega_{n,m} + h) \delta_{n,m} \right)$$

with  $\delta_{n,m} = \mathbf{1}_{(n,m) \in \tau}$ , with  $\tau$  is a two dimensional renewal

# DNA denaturation and two-dimensional renewal

Generalized Poland Scheraga model [Garel, Orland 2004]

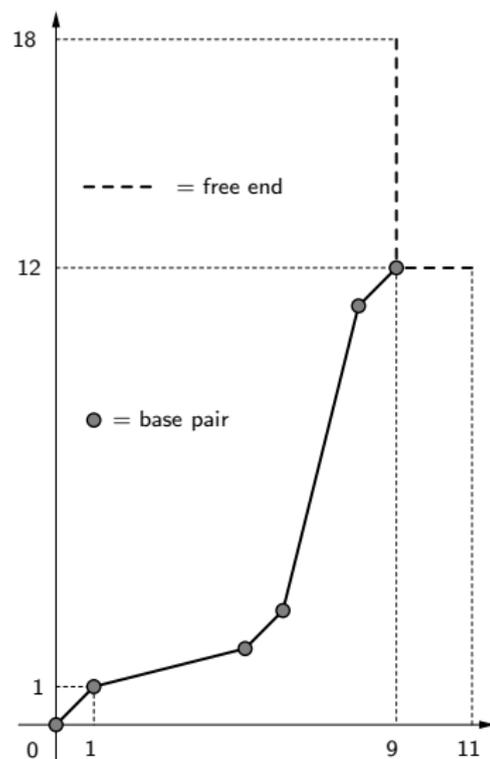


- Unequal strand length (12 and 19)
- The configuration is determined by the six base pairs (1, 1), (2, 2), (6, 3), (7, 4), (9, 12) and (10, 13).

# DNA denaturation and two-dimensional renewal

Same trajectory,  
two-dimensional renewal representation.

A  $d$ -dimensional renewal is a  $d$ -dim. walk with (componentwise) positive increments like in the one dimensional case ( $d = 1$ )



# Two-dimensional renewal pinning

Back to the definition of the model:

$$\frac{d\mathbf{P}_{N,M,\omega}}{d\mathbf{P}}(\tau) = \frac{1}{Z_{N,M,\omega}} \exp \left( \sum_{n=1}^N \sum_{m=1}^M (\beta\omega_{n,m} + h) \delta_{n,m} \right)$$

with  $\delta_{n,m} = \mathbf{1}_{(n,m) \in \tau}$ , with  $\tau$  is a two dimensional renewal

# Two-dimensional renewal pinning

Back to the definition of the model:

$$\frac{d\mathbf{P}_{N,M,\omega}}{d\mathbf{P}}(\tau) = \frac{1}{Z_{N,M,\omega}} \exp \left( \sum_{n=1}^N \sum_{m=1}^M (\beta\omega_{n,m} + h) \delta_{n,m} \right)$$

with  $\delta_{n,m} = \mathbf{1}_{(n,m) \in \tau}$ , with  $\tau$  is a two dimensional renewal with inter-arrival

$$\mathbf{P}(\tau_1 = (n, m)) \sim \frac{C_K}{(n+m)^{2+\alpha}}$$

DNA modeling: math language, but it is the choice of [Garel, Orland 2004], [Neher, Gerland 2006].

# Two-dimensional renewal pinning

Back to the definition of the model:

$$\frac{d\mathbf{P}_{N,M,\omega}}{d\mathbf{P}}(\tau) = \frac{1}{Z_{N,M,\omega}} \exp \left( \sum_{n=1}^N \sum_{m=1}^M (\beta\omega_{n,m} + h) \delta_{n,m} \right)$$

with  $\delta_{n,m} = \mathbf{1}_{(n,m) \in \tau}$ , with  $\tau$  is a two dimensional renewal with inter-arrival

$$\mathbf{P}(\tau_1 = (n, m)) \sim \frac{C_K}{(n+m)^{2+\alpha}}$$

DNA modeling: math language, but it is the choice of [Garel, Orland 2004], [Neher, Gerland 2006].

The free energy density is ( $\gamma > 0$ )

$$F_\gamma(\beta, h) := \lim_{\substack{N, M \rightarrow \infty \\ M/N \sim \gamma}} \frac{1}{N} \mathbb{E} \log Z_{N,M,\omega}$$

and the localization transition is again between  $F_\gamma(\beta, h) = 0$  and  $F_\gamma(\beta, h) > 0$ .

# Two-dimensional renewal pinning

Again (a priori somewhat surprising), homogeneous ( $\beta = 0$ ) exactly solvable, but richer behavior!

[G., Khatib 2016], [Berger, G., Khatib]

# Two-dimensional renewal pinning

Again (a priori somewhat surprising), homogeneous ( $\beta = 0$ ) exactly solvable, but richer behavior!

[G., Khatib 2016], [Berger, G., Khatib]

Summary of results:

- Localization transition and determination of  $h_c(0)$

# Two-dimensional renewal pinning

Again (a priori somewhat surprising), homogeneous ( $\beta = 0$ ) exactly solvable, but richer behavior!

[G., Khatib 2016], [Berger, G., Khatib]

Summary of results:

- Localization transition and determination of  $h_c(0)$
- Critical behavior at the localization transition: critical exponent is still  $\max(1, 1/\alpha)$

# Two-dimensional renewal pinning

Again (a priori somewhat surprising), homogeneous ( $\beta = 0$ ) exactly solvable, but richer behavior!

[G., Khatib 2016], [Berger, G., Khatib]

Summary of results:

- Localization transition and determination of  $h_c(0)$
- Critical behavior at the localization transition: critical exponent is still  $\max(1, 1/\alpha)$
- Other non analyticities (i.e., phase transitions) when  $F(0, h) > 0$

# Two-dimensional renewal pinning

Again (a priori somewhat surprising), homogeneous ( $\beta = 0$ ) exactly solvable, but richer behavior!

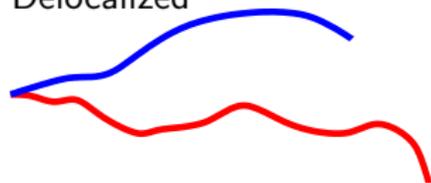
[G., Khatib 2016], [Berger, G., Khatib]

Summary of results:

- Localization transition and determination of  $h_c(0)$
- Critical behavior at the localization transition: critical exponent is still  $\max(1, 1/\alpha)$
- Other non analyticities (i.e., phase transitions) when  $F(0, h) > 0$
- Harris criterion program for IID disorder: this time disorder is irrelevant for  $\alpha < 1$ , i.e.  $h_c(\beta) = h_c(0) + \beta^2/2$  for  $\beta < \beta_0$  and  $\nu_q = \nu_a$ , and relevant for  $\alpha > 1$ , i.e.  $h_c(\beta) < h_c(0) + \beta^2/2$  for  $\beta > 0$

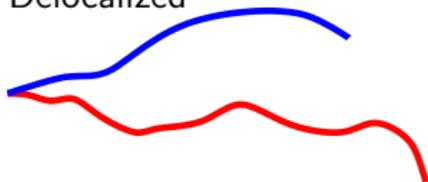
# Two-dimensional renewal pinning: paths and transitions

Delocalized

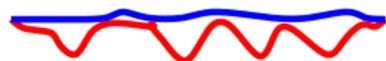


# Two-dimensional renewal pinning: paths and transitions

Delocalized



Localized



Localized



Localized



# Sum-up

- I have given an overview of the "state of the art" for pinning models with emphasis on the role and effect of disorder

# Sum-up

- I have given an overview of the "state of the art" for pinning models with emphasis on the role and effect of disorder
- Remarkably advanced understanding of the effect of disorder in this class of models

# Sum-up

- I have given an overview of the "state of the art" for pinning models with emphasis on the role and effect of disorder
- Remarkably advanced understanding of the effect of disorder in this class of models
- Still, one major open issue is unsolved: when disorder is relevant, what is the critical behavior of the quenched system?

# Sum-up

- I have given an overview of the "state of the art" for pinning models with emphasis on the role and effect of disorder
- Remarkably advanced understanding of the effect of disorder in this class of models
- Still, one major open issue is unsolved: when disorder is relevant, what is the critical behavior of the quenched system?
- Another major open issue: making the Harris criterion program rigorous for other classes of statistical mechanics systems